

Between-Group Metrics

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SUMMARY

In canonical analysis with more variables than samples, it is shown that, as well as the usual canonical means in the range-space of the within-groups dispersion matrix, canonical means may be defined in its null space. In the range space we have the usual Mahalanobis metric; in the null space explicit expressions are given and interpreted for a new metric.

Keywords: Between-group distances, Canonical analysis, Mahalanobis distance

1. INTRODUCTION

In Canonical Variate Analysis measurements on each of p variables for n samples are distributed among k groups of sizes $n_1 + n_2 + \dots + n_k = n$. These measurements are available in an $n \times p$ matrix X , assumed column-centered, and therefore of rank at most $\min(n - 1, p)$, with group-membership given in an $n \times k$ indicator matrix G . Here, $g_{ij} = 1$ when the i th sample belongs to the j th group but otherwise G is zero. Thus $G1 = 1$ and $1'G = 1'N$, where $N = \text{diag}(n_1, n_2, \dots, n_k) = G'G$; we also write ${}_nH_n = {}_nG_k N_k^{-1} G'_n$.

49 With this notation, the usual between and within-group orthogonal decomposition:

$$50 \quad {}_nX_p = {}_nG_k N_k^{-1} G'_n X_p + [I - {}_nG_k N_k^{-1} G'_n] {}_nX_p = HX + (I - H)X$$

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52 has an associated analysis of variance

$$53 \quad {}_pX'_n X_p = {}_pX'_n H_n X_p + {}_pX'_n (I - H)_n X_p$$

54 expressing that the Total sum-of-squares (T) is the sum of the Between-Group sum-of-
55 squares (B) and the Within-Group sum-of-squares (W). Note that the n rows of HX
56 repeat the k different means n_1, n_2, \dots, n_k times; to get each mean only once, we require
57 $N^{-1}GX$ which we write as \bar{X} .

58 In classical canonical variate analysis, the spectral decomposition $W = U\Sigma^2U'$ un-
59 derpins the transformation to canonical variables XL where $L = U\Sigma^{-1}$. These define
60 canonical means HXL with inner-products $(HXL)(HXL)' = HXW^{-1}X'H$ that use
61 the metric $LL' = W^{-1}$ to generate Mahalanobis distances between the canonical means;
62 note that $L'WL = I$. The rank of the canonical means is $k - 1$ (or less) but they may be
63 approximated in a smaller space, by using a conventional principal components analysis.
64 These two steps (i) define a metric, followed by (ii) a principal components analysis, are
65 usually subsumed into a single two-sided eigenvalue calculation but the two-step process
66 is better for understanding the following.

67 The above requires that W has full rank p . The case when $p > n$ is increasingly im-
68 portant where much of the interest is in overcoming computational difficulties, perhaps
69 reducing the number of variables by identifying and rejecting those deemed irrelevant
70 or by focussing on some form of functional multivariate analysis (see e.g. Krzanowski,
71 1995; Mertens, 1998). Here, we explore a novel structural property of canonical analysis
72 that occurs when $p > n$. When $p > n$ then $\text{rank}(W) = n - k$ and $\text{rank}(T) = n - 1$ and
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97 W does not have an ordinary inverse so the Mahalanobis metric is undefined. This need
 98 not be a major problem, because we may express the spectral decomposition

$$99 \quad W = (U_1, U_0) \begin{pmatrix} \Sigma^2 & \\ & 0 \end{pmatrix} \begin{pmatrix} U_1' \\ U_0' \end{pmatrix}$$

100 where U_1 are the eigenvectors in the range space of W and U_0 those in its null space.

101 Then, we may define canonical means HXL where now $L = U_1 \Sigma^{-1}$ in the range space.

102 No longer is $L'WL = I$ but rather $L'WL = I_{n-k}$. Then $(L'WL)(L'WL) = I_{n-k} = L'WL$.

103 We may write this:

$$104 \quad \begin{pmatrix} L' \\ U_0' \end{pmatrix} WLL'W(L, U_0) = \begin{pmatrix} L' \\ U_0' \end{pmatrix} W(L, U_0)$$

105 which, because (L, U_0) is non-singular, gives $W(LL')W = W$ showing that the metric

106 is now a generalised inverse, rather than an inverse, of W . With this minor change, we

107 may proceed as before with a principal components analysis. An interesting thing is that

108 canonical means may also be defined in the null space. This follows from noting that the

109 null vectors satisfy:

$$110 \quad X'(I - H)XU_0 = 0$$

111 and so

$$112 \quad XU_0 = HXU_0. \tag{1}$$

113 Note that the k different means are repeated n_1, n_2, \dots, n_k times in the n rows of both

114 XU_0 and, equivalently, HXU_0 . Being null vectors of W , the canonical variables XU_0

115 have zero variability within groups, but the corresponding canonical means HXU_0 have

116 non-zero sums-of-squares. Evidently, the computation of HXU_0 is straightforward, as is

117 any subsequent principal components analysis; an example is given by Gower & Albers

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145 ('Canonical Analysis: Ranks, Ratios and Fits', in preparation). For a fuller understanding
146 it is interesting to ask what functional form, analogous to Mahalanobis distance in the
147 range space, is taken by the distance d_{ij} between the i th and j th canonical means in the
148 null space of W . This is our main objective below but first we have to address a minor
149 but troublesome technical matter.

150 The total dispersion $T = X'X$ has rank $n - 1$ so implying an extensive null space
151 of rank $p - n + 1$; this null space is also common to the null spaces of B and W . This
152 common null space is uninteresting; we are concerned only with the additional null spaces
153 of W and B that are in the range space of T , especially the intersection of the range space
154 of T and the null space of W which normally has dimension/rank $k - 1$. To simplify the
155 following development we assume that the common null space has been eliminated by
156 taking the spectral decomposition $T = V\Lambda V'$ and redefining X as XV . Throughout the
157 following, we assume that X has been so redefined.

158 This initialisation to give X with $n - 1$ columns, eliminates the common null space
159 from the dispersion matrices T , B and W . However, it does not remove null items from
160 X itself. Indeed the vector 1 , which eliminates the general mean, is one such null vector
161 and is what gives rise to the rather extensive algebraic manipulations required in the
162 following. Linear combinations among the rows of X will generate additional null vectors
163 in the common null space. The position is complicated, because such linear combinations
164 may be of two, not mutually exclusive, kinds (i) linear combinations within groups and
165 (ii) linear combinations among the group means. Loss of rank within groups merely
166 reduces the number of columns of the redefined X but to handle all variants that include
167 (ii) is not trivial and would greatly extend this short paper. Therefore, apart from some
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193 passing concluding remarks, throughout the following, we assume that $\text{rank}\bar{X} = k - 1$
 194 and that $\text{rank}X \leq n - 1$.

197 2. DERIVATION OF d_{ij}^2

198 From here on we shall be working in the null space of W so we drop the suffix from
 199 U_0 . Starting from $XU = HXU$ for the null-vectors of W , as in (1), we have that

$$\begin{aligned} 201 \quad {}_nX_{n-1}U_{k-1} &= G(G'G)^{-1}G'XU \\ 202 &= {}_nG_kA_{k-1} \end{aligned}$$

203 where ${}_kA_{k-1} = \bar{X}U$ are the k group-mean coordinates given in repeated form in XU .

204 Thus, the calculation $A = (G'G)^{-1}G'XU$ is an expression whose rows give coordi-
 205 nates that generate the distance d_{ij} between each pair of group-mean coordinates. We
 206 need an explicit expression for d_{ij}^2 . We do not require A itself, which has the usual
 207 rotational indeterminacy, but only AA' . Then, $d_{ij}^2 = (AA')_{ii} + (AA')_{jj} - 2(AA')_{ij}$. Be-
 208 cause X is centred, $1'X = 0$ and so $1'XU = 1'GA = 1'NA = 0$. Also, $U'U = I_{k-1}$. From
 209 (1) $U'U = ((X'X)^{-1}X'GA)'((X'X)^{-1}X'GA) = A'PA$ where $P = G'X(\Lambda^{-2})X'G$ with
 210 $X'X = \Lambda$. We have $1'P = 1'G'X(\Lambda^{-2})X'G = 1'X(\Lambda^{-2})X'G = 0$. So we have to solve
 211 $A'PA = I$ for AA' . The difficulty is that both A and P are singular (with rank $k - 1$).

212 Hence, consider

$$\begin{pmatrix} A'N \\ \frac{1}{n}1'N \end{pmatrix} (Q + \lambda 11')(NA, \frac{1}{n}N1) \quad (2)$$

241 where $Q = N^{-1}PN^{-1} = \bar{X}\Lambda^{-2}\bar{X}'$. The introduction of λ may seem arbitrary but we
 242 shall show that it has no substantive effect. On expansion, (2) becomes $\begin{pmatrix} I & 0 \\ 0 & \lambda \end{pmatrix}$ giving:
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$$244 \quad Q + \lambda 11' = \begin{pmatrix} A'N \\ \frac{1}{n}1'N \end{pmatrix}^{-1} \begin{pmatrix} I \\ \lambda \end{pmatrix} (NA, \frac{1}{n}N1)^{-1}$$

246 and

$$247 \quad (Q + \lambda 11')^{-1} = (NA, \frac{1}{n}N1) \begin{pmatrix} I \\ \frac{1}{\lambda} \end{pmatrix} \begin{pmatrix} A'N \\ \frac{1}{n}1'N \end{pmatrix}.$$

249 Thus,

$$250 \quad N^{-1} (Q + \lambda 11')^{-1} N^{-1} = \frac{AA' + 11'}{\lambda n^2}. \quad (3)$$

251 From 3 we may calculate d_{ij}^2 . The constant term $11'/\lambda n^2$ has no effect on derived distances
 252 and we shall show that $(Q + \lambda 11')^{-1}$ also is invariant to non-zero choices of λ . Thus (3)
 253 contains everything needed for finding AA' but the evaluation of $(Q + \lambda 11')^{-1}$ needs
 254 some care, because Q is singular and the equivalent of (A4) is unavailable.

255 For simplicity, we derive d_{12}^2 , the other values of d_{ij}^2 following by symmetry. Notation
 256 is established via

$$257 \quad C = \left(\begin{array}{c|c|c} c_{11} & c_{21} & c'_1 \\ \hline c_{12} & c_{22} & c'_2 \\ \hline c_1 & c_2 & C_{12} \end{array} \right) \quad \text{where} \quad \left\{ \begin{array}{l} C = Q + \lambda 11' \\ C_{12} = Q_{12} + \lambda 11' \\ c_1 = q_1 + \lambda 1 \\ c_2 = q_2 + \lambda 1 \\ (c_{11}, c_{12}, c_{22}) = (q_{11}, q_{12}, q_{22}) + \lambda \end{array} \right.$$

262 From (3)

$$263 \quad \Delta d_{12}^2 = \frac{1}{n_1^2} \det \begin{pmatrix} c_{22} & c'_2 \\ c_2 & C_{12} \end{pmatrix} + \frac{1}{n_2^2} \det \begin{pmatrix} c_{11} & c'_1 \\ c_1 & C_{12} \end{pmatrix} + \frac{2}{n_1 n_2} \det \begin{pmatrix} c_{12} & c'_2 \\ c_1 & C_{12} \end{pmatrix}, \quad (4)$$

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289 where $\Delta = \det C$, and the determinants are the cofactors of c_{11} , c_{22} and c_{12} . Using (A1),
 290 (4) becomes

$$291 \Delta d_{12}^2 = \det C_{12} \left[\frac{1}{n_1^2} (c_{22} - c_2' C_{12}^{-1} c_2)^{-1} + \frac{1}{n_2^2} (c_{11} - c_1' C_{12}^{-1} c_1)^{-1} + \frac{2}{n_1 n_2} (c_{12} - c_1' C_{12}^{-1} c_2)^{-1} \right],$$

292 which simplifies to

$$293 n_1^2 n_2^2 \Delta d_{12}^2 = \det C_{12} \left[n_1^2 c_{11} + n_2^2 c_{22} + 2n_1 n_2 c_{12} - (n_1 c_1 + n_2 c_2)' C_{12}^{-1} (n_1 c_1 + n_2 c_2) \right].$$

$$295 \tag{5}$$

296 Using (A1) and (A2) we have

$$297 (n_1 c_1 + n_2 c_2)' C_{12}^{-1} (n_1 c_1 + n_2 c_2) = (C_{12} M 1 - \lambda n 1)' C_{12}^{-1} (C_{12} M 1 - \lambda n 1)$$

$$298 = 1' M C_{12} M 1 - 2\lambda n 1' M 1 + \lambda^2 n^2 1' C_{12}^{-1} 1. \tag{6}$$

299 Bringing everything together using (6), (A4) and (A5), and expanding in terms of q_{ij} ,
 300 (5) becomes

$$301 n_1^2 n_2^2 \Delta d_{12}^2 = (1 + \lambda 1' Q_{12} 1) \det Q_{12} \left[n_1^2 q_{11} + n_2^2 q_{22} + 2n_1 n_2 q_{12} + \lambda (n_1 + n_2)^2 \right.$$

$$302 \left. - 1' M Q_{12} M 1 - \lambda (1' M 1)^2 + 2\lambda n 1' M 1 - \lambda^2 n^2 \left(Q_{12}^{-1} - \frac{\lambda Q_{12}^{-1} 1 1' Q_{12}^{-1}}{1 + \lambda 1' Q_{12}^{-1} 1} \right) \right]$$

303 which, on using (A1), simplifies to

$$304 n_1^2 n_2^2 \Delta d_{12}^2 = \det Q_{12} \left[\lambda (n_1 + n_2)^2 + \lambda (n - n_1 - n_2)(n + n_1 + n_2) - \frac{\lambda n^2 1' Q_{12}^{-1} 1}{1 + \lambda 1' Q_{12}^{-1} 1} \right] (1 + \lambda 1' Q_{12}^{-1} 1)$$

$$305 = \det Q_{12} \left[\lambda n^2 (1 + \lambda 1' Q_{12}^{-1} 1) - \lambda^2 n^2 1' Q_{12}^{-1} 1 \right]$$

$$306 = \lambda n^2 \det Q_{12}. \tag{7}$$

307 This simple result shows that d_{ij}^2 is proportional to $\det Q_{ij}$.

308 To show that d_{ij}^2 is independent of λ requires an analysis of $\Delta = \det(Q + \lambda 11')$. Let
 309 $R = Q + 11'$. Then $I = R^{-1}Q + R^{-1}11'$ and so

$$310 1' I N 1 = 1' (R^{-1}Q + R^{-1}11') N 1$$

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$$1'N1 = (1'R^{-1}1)(1'N1)$$

$$1'R^{-1}1 = 1.$$

Then

$$\begin{aligned} \Delta &= \det(R + (\lambda - 1)11') \\ &= \det R (1 + (\lambda - 1)1'R^{-1}1) \\ &= \lambda \det R \\ &= \lambda \det(Q + 11'). \end{aligned}$$

That $R \neq 0$ is guaranteed by our assumption that $\bar{X} = k - 1$ made at the end of Section 1.

Thus, finally, (7) becomes

$$d_{12}^2 = \frac{1}{\det(Q + 11')} \frac{n^2}{n_1^2 n_2^2} \det Q_{12} \quad (8)$$

showing that d_{ij}^2 depends only on the group sizes and $\det Q_{12}$. Recall that $Q = \bar{X}\Lambda^{-1}\bar{X}'$ and that Q_{12} is obtained from Q by striking out its first two rows and columns.

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3. ALTERNATIVE EXPRESSIONS

The expression for Δ has many forms. Of special interest is that derived from writing

$$\Delta = \det \left(\begin{array}{cc|c} c_{11} & c_{12} & c'_1 \\ c_{12} & c_{22} & c'_2 \\ \hline c_1 & c_2 & C_{12} \end{array} \right).$$

Multiplying the first row by n_1 and then adding q_i ($i = 2, \dots, k$) times the other rows, replaces the first row by $n\lambda 1$. This shows that λ may be subtracted from rows $2, \dots, k$

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$$n_1 \Delta = \lambda n \det \left(\begin{array}{cc|c} 1 & 1 & 1' \\ q_{12} & q_{22} & q'_2 \\ \hline q_1 & q_2 & Q_{12} \end{array} \right)$$

390 whence similar operations on the columns give

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$$\begin{aligned} n_1^2 \Delta &= \lambda n^2 \det \left(\begin{array}{c|cc} 1 & 1 & 1' \\ 0 & q_{22} & q'_2 \\ \hline 0 & q_2 & Q_{12} \end{array} \right) \\ &= \lambda n^2 (q_{22} - q'_2 Q_{12}^{-1} q_2) \det Q_{12}. \end{aligned} \quad (9)$$

395 Similar expressions may be derived by annihilating the second row/column and the first

396 row and second column to give

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$$\left. \begin{aligned} n_2^2 \Delta &= \lambda n^2 (q_{11} - q'_1 Q_{12}^{-1} q_1) \det Q_{12} \\ n_1^2 \Delta &= \lambda n^2 (q_{22} - q'_2 Q_{12}^{-1} q_2) \det Q_{12} \\ -n_1 n_2 \Delta &= \lambda n^2 (q_{12} - q'_1 Q_{12}^{-1} q_2) \det Q_{12} \end{aligned} \right\}. \quad (10)$$

401 Combining, gives the symmetric form

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$$(n_1 + n_2)^2 \Delta = \lambda n^2 [(q_{11} + q_{22} - 2q_{12}) - (q_1 - q_2)' Q_{12}^{-1} (q_1 - q_2)] \det Q_{12}$$

404 which, on substitution into (8) gives

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$$d_{12}^2 = \frac{(n_1 + n_2)^2}{n_1^2 n_2^2} [(\bar{x}_1 - \bar{x}_2)' \Lambda^{-2} (\bar{x}_1 - \bar{x}_2) - (q_1 - q_2)' Q_{12}^{-1} (q_1 - q_2)]^{-1}. \quad (11)$$

407 Other substitutions for Δ given by (10) give alternative, less symmetric, expressions for

408 d_{ij}^2 .

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4. INTERPRETATION

From (11), for $k = 2$ we immediately have

$$d_{12}^{-2} = \frac{n_1^2 n_2^2}{n^2} (\bar{x}_1 - \bar{x}_2)' \Lambda^{-2} (\bar{x}_1 - \bar{x}_2). \quad (12)$$

We next examine the part of expression (11) that is enclosed in square brackets. We have

$$\bar{X} = \begin{pmatrix} \bar{x}_1 \\ \bar{x}_2 \\ \hline \bar{X}_{12} \end{pmatrix} \quad \text{and} \quad \begin{aligned} (q_1 - q_2) &= (\bar{x}_1 - \bar{x}_2) \Lambda^{-2} \bar{X}_{12}' \\ Q_{12} &= \bar{X}_{12} \Lambda^{-2} \bar{X}_{12}' \end{aligned}.$$

Hence,

$$d_{12}^{-2} = \frac{n_1^2 n_2^2}{(n_1 + n_2)^2} \left[(\bar{x}_1 - \bar{x}_2)' \Lambda^{-2} (\bar{x}_1 - \bar{x}_2) - (\bar{x}_1 - \bar{x}_2)' \Lambda^{-1} R \Lambda^{-1} (\bar{x}_1 - \bar{x}_2) \right]$$

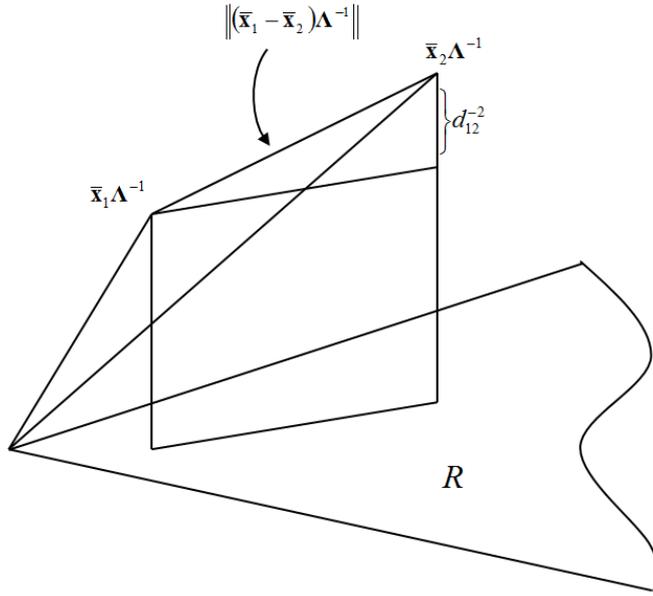
where $R = \Lambda^{-1} \bar{X}_{12} (\bar{X}_{12} \Lambda^{-2} \bar{X}_{12})^{-1} \bar{X}_{12} \Lambda^{-1}$ represents projection in the metric Λ onto the space spanned by the $k - 2$ rows of $\bar{X}_{12} \Lambda$. Thus

$$d_{12}^{-2} = \frac{n_1^2 n_2^2}{(n_1 + n_2)^2} \left[(\bar{x}_1 - \bar{x}_2)' \Lambda^{-1} (I - R) \Lambda^{-1} (\bar{x}_1 - \bar{x}_2) \right]. \quad (13)$$

Thus, expressions (6), (11) and (13) are our main results.

The interpretation of (13), visualised in Figure 1, is that d_{12}^{-1} is a measure of how far the space of \bar{x}_1, \bar{x}_2 is from the space spanned by $\bar{x}_3, \dots, \bar{x}_k$. These results are expressed in terms of the Λ -metric but this is a function of the initial rescaling and vanishes on transforming back to the scales of the original variables.

If \bar{x}_1 and \bar{x}_2 lie in R then (13) gives a null projection onto $I - R$ but this case is excluded by our assumption that $\text{rank} \bar{X} = k - 1$. An indication of the corresponding results when $\text{rank} \bar{X} < k - 1$ is given by the collinearity case $k = 3$ and $\text{rank} \bar{X} = 1$. Then, without giving the detailed derivation, it can be shown that d_{12}^2 is proportional to $(\bar{x}_1 - \bar{x}_2)' \Lambda^{-2} (\bar{x}_1 - \bar{x}_2)$. That the projection term vanishes is consistent and plausible but we do not know whether it generalises to other reduced rank cases.



490 Fig. 1. R is the space spanned by the means
 491 $\bar{x}_3, \dots, \bar{x}_k$. The inverse of the distance between
 492 groups 1 and 2 in the intersection space is given by
 493 the illustrated projection onto the space orthogonal
 494 to R .

495 5. BASIC FORMULAE

496 This appendix contains some basic results for convenience of reference. Most are well-known
 497 and, apart from (A1) and (A2), no derivations are supplied.

498 Using $1'NQ = 0$ shows that

$$499 n_1q_{11} + n_2q_{12} + 1'Mq_1 = 0$$

$$500 n_1q_{12} + n_2q_{22} + 1'Mq_2 = 0$$

$$501 n_1q_1 + n_2q_2 + Q_{12}M1 = 0$$

502 where $N = \text{diag}(n_1, n_2, M)$. It follows that

$$503 n_1^2q_{11} + n_2^2q_{22} + 2n_1n_2q_{12} = -1'M(n_1q_1 + n_2q_2) = 1'MQ_{12}M1. \quad (\text{A1})$$

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529 Similarly,

$$\begin{aligned}
 530 \quad n_1 c_1 + n_2 c_2 &= n_1 q_1 + n_2 q_2 + \lambda(n_1 + n_2)1 \\
 531 \quad &= -Q_{12}M1 + \lambda(n_1 + n_2)1 \\
 532 \quad &= -C_{12}M1 + \lambda n1. \tag{A2}
 \end{aligned}$$

533 Furthermore, we need

$$534 \quad \det \begin{pmatrix} \alpha & a' \\ a & A \end{pmatrix} = (\alpha - a' A^{-1} a) \det A, \tag{A3}$$

$$536 \quad \det C_{12} = \det(Q_{12} + \lambda 11') = \det Q_{12} (1 + \lambda 1' Q_{12}^{-1} 1), \tag{A4}$$

537 and

$$538 \quad C_{12}^{-1} = (Q_{12} + \lambda 11')^{-1} = Q_{12}^{-1} - \frac{\lambda Q_{12}^{-1} 11' Q_{12}^{-1}}{1 + \lambda 1' Q_{12}^{-1} 1}. \tag{A5}$$

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